







MECHANICS OF 2D MATERIALS



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Outline

Stretching

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Stress Strain curve
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Mechanical Properties

Young's modulus

Strength

Ultimate strain

Toughness modulus

Size effects on energy dissipated

Linear Elastic Fracture Mechanics (LEFM)

Stress-intensity factor

Energy release rate

Fracture toughness

Size-effect on fracture strength

Quantized Fracture Mechanics (QFM)

Strength of graphene (and related materials)

Bending

Flexibility

Bending stiffness

Elastic line equation

Elastic plate equation

Exercises by me

- 1. Apply QFM for deriving the strength of realistic thus defective graphene (and related 2D materials);
- Apply LEFM for deriving the peeling force of graphene (and related 2D materials).

Exercises by you

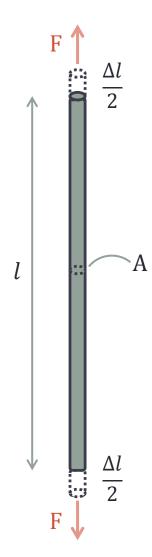
- Apply LEFM for measuring the fracture toughness of a sheet of paper;
- Apply the Bending theory for calculating the maximal curvature before fracture.

Stretching

Fiber under tension

Stress:
$$\sigma = \frac{F}{A}$$

Strain:
$$\varepsilon = \frac{\Delta l}{l}$$



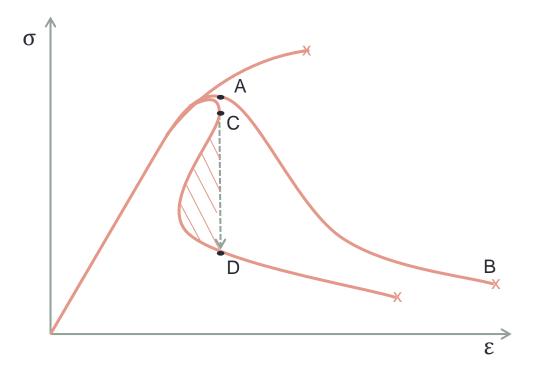
Stress-strain curve:

Signature of the material and main tool for deriving its mechanical properties

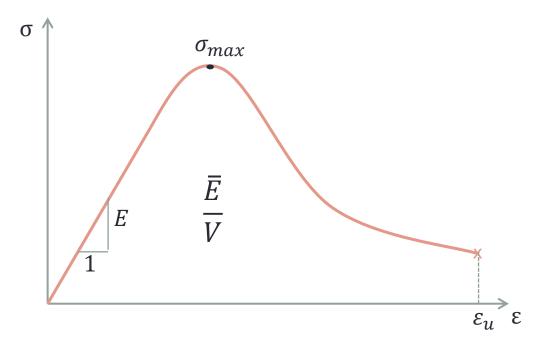
x = failure

If the curve is monotonic, a force-control is sufficient.

You can derive \widehat{AB} in displacement control, \widehat{CD} in crack-opening. The dashed area is the kinetic energy released per unit volume under displacement control.



Mechanical properties



$$\frac{d\sigma}{d\varepsilon}\Big|_0 \equiv E \equiv \text{Young's modulus}$$

 $\sigma_{max} = \text{maximum stress} \equiv \text{Strength}$

 ε_u = ultimate strain

$$\int_0^{\varepsilon_u} \sigma d\varepsilon = \frac{\bar{E}}{V}$$
 = Energy dissipated per unit volume or Toughness modulus.

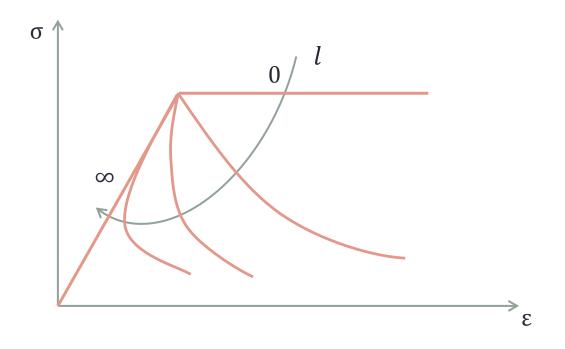
$$\bar{E} = \int F dl \qquad V = Al$$

$$\frac{\bar{E}}{V} = \int \frac{F}{A} \frac{dl}{l} = \int \sigma d\varepsilon$$

The toughness modulus is a material property only for ductile materials.

Size-effects

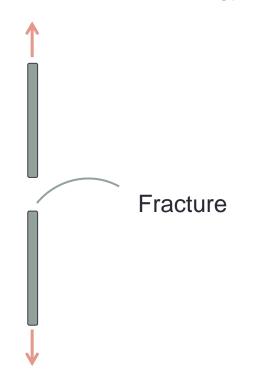
The post-critical behaviour is size-dependent especially for brittle materials.



 $l \rightarrow 0$ ductile

 $l \rightarrow \infty$ brittle

For brittle materials $\frac{\bar{E}}{V}$ has no meaning: instead of the toughness modulus we have to use the fracture energy.



$$\bar{E} = G_c A$$

 $G_c \equiv$ fracture energy or energy dissipated per unit area

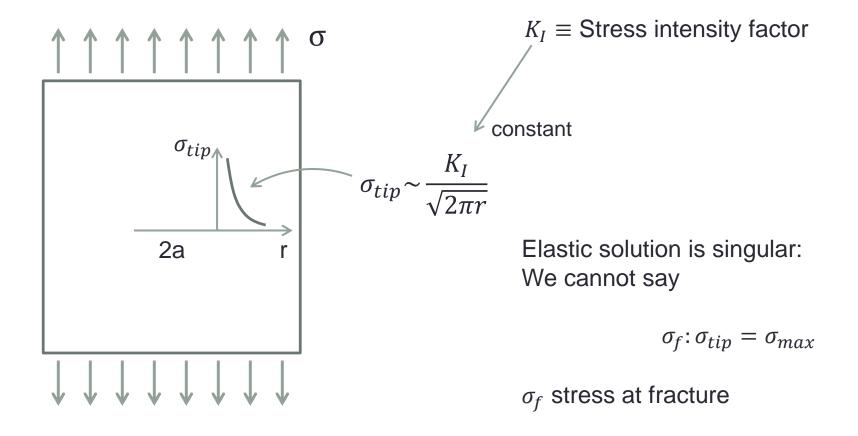
$$\int \sigma d\varepsilon \equiv \frac{\bar{E}}{V} = \frac{G_c A}{lA} = \frac{G_c}{l}$$

The reality is in between: energy dissipated on a fractal domain:

$$\frac{\overline{E}}{V} \propto \tilde{l}^{D-3} \qquad \qquad \tilde{l} = \sqrt[3]{V}$$

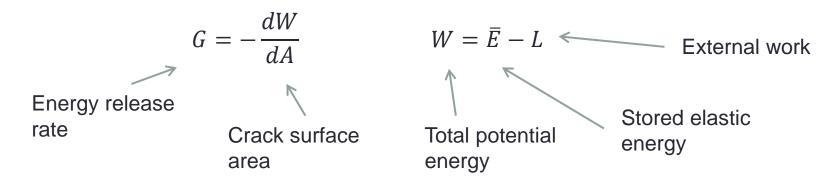
D is the fractal exponent, $2 \le D \le 3$

Fracture Mechanics



Linear Elastic Fracture Mechanics (LEFM)

Griffith's approach



Crack propagation criterion:

$$G = G_C$$

$$G_c \equiv \text{fracture energy}$$

$$G = \frac{K_I}{F}$$
 Elastic solution

 K_I only for a function of external load and geometry

$$K_I = K_{IC}$$
 $K_{IC} = \sqrt{G_C E}$ Fracture toughness

 K_I values reported in stress-intensity factor Handbooks

E.g., infinite plate (i.e. width >> crack length ~ graphene)

$$K_I = \sigma \sqrt{\pi a}$$

Strength of graphene with a crack of length 2a

$$K_I = \sigma \sqrt{\pi a} = K_{IC}$$

$$\sigma_f = \frac{K_{IC}}{\sqrt{\pi a}}$$

Assuming statistically $a \propto \tilde{l} \equiv$ structural size:

$$\sigma_f \propto \tilde{l}^{-1/2}$$

Size effect on a fracture strength



larger is weaker



problem of the scaling up....

Paradox $\sigma_f \to \infty$ $a \to 0$

With graphene pioneer Rod Ruoff we invented Quantized Fracture Mechanics (2004).

The hypothesis of the continuous crack growth is removed: existence of fracture quanta due to the discrete nature of matter.

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Quantized fracture mechanics

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ABSTRACT

A new energy-based theory, quantized fracture mechanics (QFM), is presented that modifies continuum-based fracture mechanics; stress- and strainbased QFM analogs are also proposed. The differentials in Griffith's criterion

Related papers:

N. M. Pugno, "Dynamic quantized fracture mechanics", Int. J. Fract, 2006

N. M. Pugno, "New quantized failure criteria: application to nanotubes and nanowires", Int. J. Fract, 2006

Quantized Fracture Mechanics (QFM)

$$G^* = -\frac{\Delta W}{\Delta A}$$
Quantized energy

release rate

 ΔA = fracture quantum of surface area = qt



Fracture Plate quantum thickness of length

$$G = -\frac{dW}{dA} \qquad \Delta W = -\int G dA = -\int \frac{K_I^2}{E} dA$$

$$G^* = -\frac{\Delta W}{\Delta A} = \frac{\int \frac{K_I^2}{E} dA}{\Delta A} \qquad G^* = G_c \qquad \frac{1}{\Delta A} \int \frac{K_I^2}{E} dA = \frac{K_{IC}^2}{E}$$

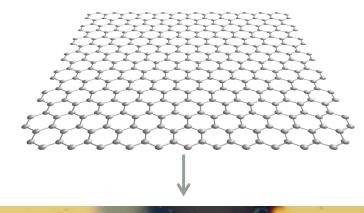
Quantized stress-intensity factor (generalized):

$$K_I^* = \sqrt{\frac{1}{\Delta A} \int K_I^2 dA} = K_{IC}$$

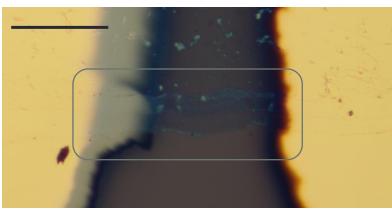
Monolayer graphene and examples

C. Lee, X. Wei, J.W. Kysar, J. Hone,

"Measurement of the elastic properties and intrinsic strength of monolayer graphene", Science, 2008:



Young's modulus $\equiv E=1$ TPa Intrinsic strength $\sigma_{int}=130$ GPa Ultimate strain $\varepsilon_u=25$ %







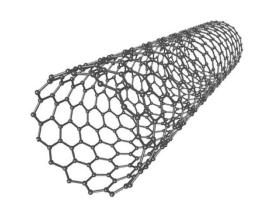
Monolayer graphene hanging on a silicon substrate (scale bar: 50µm)

Tensile test on macro samples of graphene composites

Tennis racket made of graphene (*Head* ©)

Carbon nanotubes

Min-Feng Yu, Oleg Lourie, Mark J. Dyer, Katerina Moloni, Thomas F. Kelly, Rodney S. Ruoff, "Strength and Breaking Mechanism of Multiwalled Carbon Nanotubes Under Tensile Load", Science, 2000:

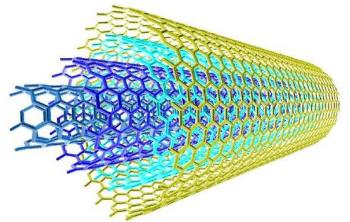


The singlewalled carbon nanotubes (SWCNTs) presents:

Young's modulus $\equiv E = 1000 \text{ GPa}$

Tensile strength $\sigma = 300 \text{ GPa}$

Ultimate strain $\varepsilon_u = 30 \%$



The MWCNTs breaks in the outermost layer ("sword-in-sheath" failure),

Young's modulus $\equiv E = 250$ to 950 GPa

Tensile strength $\sigma = 11$ to 63 GPa

Ultimate strain $\varepsilon_u = 12 \%$

Multiwalled carbon nanotubes (MWCNTs)

Exercise 1

E.g., infinite plate

$$K_I = \sigma \sqrt{\pi a}$$

$$K_{I}^{*} = \sqrt{\frac{1}{qt} \int_{a}^{a+q} \sigma^{2} \pi at \ da} = \sigma \sqrt{\frac{\pi}{2q} [(a+q)^{2} - a^{2}]} = \sigma \sqrt{\pi \left(a + \frac{q}{2}\right)}$$

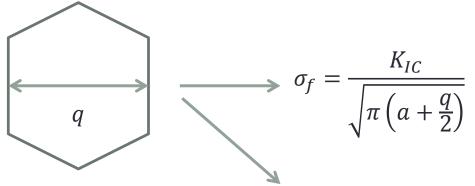
Strength of graphene $K_I^* = K_{IC}$

$$\sigma_f = \frac{K_{IC}}{\sqrt{\pi \left(a + \frac{q}{2}\right)}} \qquad a \uparrow \quad \sigma_f \downarrow \quad \text{Unstable crack growth}$$

$$q \rightarrow 0$$
 Correspondence Principle $QFM \rightarrow LEFM$

$$a \rightarrow 0$$
 $\sigma_f = \sigma_{ideal}$

$$\sigma_{ideal} = \frac{K_{IC}}{\sqrt{\pi \frac{q}{2}}} \implies q = \frac{2}{\pi} \frac{K_{IC}^2}{\sigma_{ideal}^2}$$



Very close to predictions by MD and DFT

$$\sigma_{ideal} \sim 100 \; \mathrm{GPa}$$
 $\sigma_{ideal} \sim \frac{E}{10} \quad E \sim 1 \; \mathrm{TPa}$

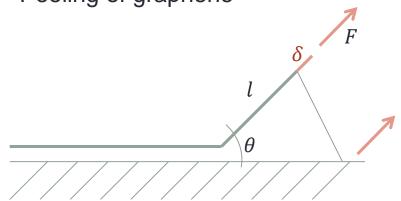
$$K_{IC} \sim 3\text{MPa}\sqrt{\text{m}}$$
(QFM 2004)

Very close to experimental measurement (2014)

P. Zhang, L. Ma, F. Fan, Z. Zeng, C. Peng, P. E. Loya, Z. Liu, Y. Gong, J. Zhang, X. Zhang, P. M. Ajayan, T. Zhu & J. Lou, *Fracture toughness of graphene*, Nature, 2014

Exercise 2





$$E \to \infty$$
 $\bar{E} \to 0$
$$\delta = l - l \cos \theta$$

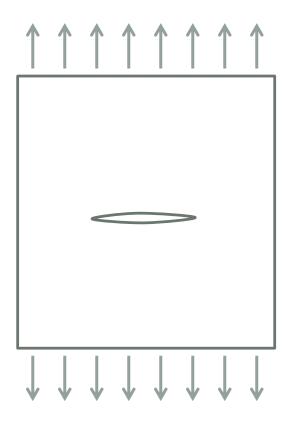
$$G = -\frac{dW}{dA} = \frac{dL}{dA} \qquad L = Fd \qquad dA = bdl$$

$$G = \frac{d}{bdl}Fl(1 - \cos\theta) = \frac{F(1 - \cos\theta)}{b}$$

$$G = G_C \implies F_C = \frac{bG_C}{1 - \cos\theta}$$

 F_C strongly dependent on θ

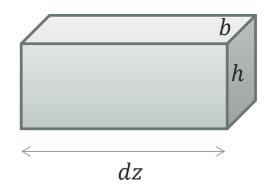
Exercise 3



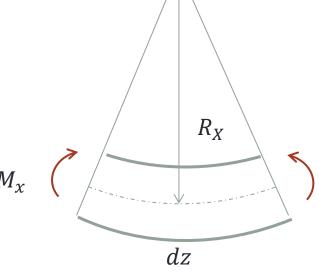
Apply LEFM for measuring the fracture toughness of a sheet of paper with a crack of a certain length in the middle:

Bending theory

Bending of beams



 $d\varphi_X$



 $d\varphi_X$

Bending Moment

 M_{χ}

Moment of inertia

$$I_{x} = \frac{bh^{3}}{12}$$

$$\sigma_{z} = \frac{M_{x}}{I_{x}}y \cong Ky \qquad M_{x} = \int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_{z}ybdy = KI_{x}$$

$$\frac{\varepsilon dz}{2} \cong y \frac{d\varphi}{2}$$

$$\chi_x = \frac{1}{R_x} = \frac{d\varphi_x}{dz} = \frac{\varepsilon_z}{y} = \frac{\sigma_z}{Ey} = \frac{M_x}{I_x E}$$

$$\chi_{max} = \frac{\varepsilon_{z,max}}{h/2}$$

Exercise 4

Derive the maximal graphene curvature.

For graphene $h = t \approx 0.34 \ nm \rightarrow \chi_{max}$ huge \rightarrow flexibility is more a structural than a material property

$$\varphi_x = -\frac{dv}{dz} \implies \frac{d^2v}{dz^2} = -\frac{M_x}{I_x E}$$

Load per unit length

$$\begin{array}{c|c}
T \uparrow & \downarrow q & \uparrow T \\
M \nearrow N & \frac{dT}{dz} = -q \\
\hline
\frac{dM}{dz} = T
\end{array}$$

$$\begin{array}{c}
\frac{d^4v}{dz^4} = \frac{q}{I_x E} \\
\text{Elastic line equation}$$

$$\frac{dT}{dz} = -q$$

$$\Rightarrow \frac{d^4v}{dz^4} = \frac{q}{I_x E}$$
Elastic line equation

Load per unit area $\nabla^4 W = \frac{q}{R}$ For plates:

Elastic plane equation

Bending stiffness
$$D = \frac{Et^3}{12(1-v^2)}$$